

ESLW_Drivers

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Complementary to the previous ESL workshops, which addressed solvers for a variety of specific representations of the KS problem, "drivers" targeted the problem of extracting existing KS solvers from major Electronic Structure codes and refactor them as libraries in a representation-agnostic fashion, resulting in stronger code interoperability.

The two main iterative eigensolvers employed in the *pw.x* code of the *Quantum ESPRESSO* distribution have been completely disentangled from the rest of the code and converted to libraries. The solvers make use of a Linear Algebra domain-specific library LAXlib, developed within the MaX CoE, which is interfaced with ELPA and ScalaPack.

Solvers exploit MPI parallelization and in addition to basis-set component distribution, a parallelization over target states is possible, as well as a specific parallelization for the dense linear algebra.

Generic k-point as well as Gamma specific versions of the solvers are included. The Reverse Communication Interface (RCI) paradigm, allowing for a complete abstraction from the basis type and the interface used to perform the matrix-vector operations, has also been implemented for one of the solvers.



Diagonalization of H_{KS} is a major step in the scf solution of any system.

In `pw.x` two methods are implemented:

- Davidson diagonalization

- efficient in terms of number of H_{psi} required
- memory intensive: requires a work space up to

$$(1 + 3 * david) * nbnd * npwx$$

and diagonalization of matrices up to

$$david * nbnd \times david * nbnd$$

where *david* is by default 4, but can be reduced to 2

- Conjugate Gradient

- memory friendly: bands are dealt with one at a time.

- the need to orthogonalize to lower states makes it intrinsically sequential and not efficient for large systems.

Davidson Diagonalization

• Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

• Eigenpairs of the reduced Hamiltonian

$$\tilde{H}_{ij} = \langle \phi_i^{(n)} | H_{KS} | \phi_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \phi_i^{(n)} | S | \phi_j^{(n)} \rangle$$

• Build the correction vectors $|\tilde{\phi}_i^{(n)}\rangle$

$$|\tilde{\phi}_i^{(n)}\rangle = (H_{diag} - \varepsilon_i S_{diag})^{-1} (H_{KS} - \varepsilon_i S) |\phi_i^{(n)}\rangle$$

• Build an extended reduced Hamiltonian

$$\tilde{H}_{ij} = \langle \phi_i^{(n)} / \tilde{\phi}_i^{(n)} | H_{KS} | \phi_j^{(n)} / \tilde{\phi}_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \phi_i^{(n)} / \tilde{\phi}_i^{(n)} | S | \phi_j^{(n)} / \tilde{\phi}_j^{(n)} \rangle$$

• Diagonalize the small $2nbnd \times 2nbnd$ reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \quad \longrightarrow \quad \{|\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)}\}$$

• Repeat if needed in order to improve the solution

$\rightarrow 3nbnd \times 3nbnd \rightarrow 4nbnd \times 4nbnd \dots \rightarrow \underline{nbnd \times nbnd}$

- Davidson diagonalization

- efficient in terms of number of H_{psi} required
- memory intensive: requires a work space up to

$$(1 + 3 * david) * nbnd * npwx$$

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where *david* is by default 4, but can be reduced to 2

- routines

- *regterg* , *cegterg* real/cmplx eigen iterative generalized

- *h_psi*, *s_psi*, *g_psi*

- *rdiaghg*, *cdiaghg* real/cmplx diagonalization H generalized

Conjugate Gradient

- For each band, given a trial eigenpair: $\{|\phi_i^{(n)}\rangle, \varepsilon_i\}$

- Minimize the single particle energy

$$E(|\phi_i\rangle) = \langle \phi_i | H_{KS} | \phi_i \rangle$$

by (pre-conditioned) CG method

subject to the constraints

$$\langle \phi_i | S | \phi_j \rangle = \delta_{ij}, \quad \forall j \leq i$$

- Repeat for next band until completed

- Conjugate gradient

- memory friendly: bands are dealt with one at a time.
- the need to orthogonalize to lower states makes it intrinsically sequential and not efficient for large systems.

- routines

- rcgdiagg , ccgdiagg *real/cmplx CG diagonalization generalize*
- rotate_wfc_gamma, rotate_wfc_k *real/cmplx initial diag*
- h_1psi, s_1psi
 - * preconditioning

The software developed during the Workshop is hosted by the e-cam gitlab server as a public sub-project of the ESL initiative (gitlab.e-cam2020/esl/ESLW_Drivers). It will be released under LGPL 2.1+

A toy code implementing the Cohen-Bergstresser empirical pseudopotential method is included to exemplify the use of the solvers and allow a quick test of their functionalities.

A git branch to contribute additional toy codes based on different basis set representations, tight-binding or real-space, to explicitly demonstrate the representation-agnostic nature of the solvers has been created and is under development.

A version of the PPCG algorithm is being ported.

Contribution to the KS_Solvers library of additional solvers from other codes is also foreseen.



https://gitlab.e-cam2020.eu/esl/ESLW_Drivers

| | |
|-----------------|---|
| CB_toy_code/Doc | <i>empty</i> |
| /examples | <i>contains inputs and ref. outputs</i> |
| /src | <i>contains simple code mains</i> |
| FFTXlib | <i>fft library used by CB_toy_code</i> |
| KS_Solvers/CG | <i>band-by-band CG</i> |
| /Davidson | <i>Davidson iterative diagonalization</i> |
| /Davidson_RCI | <i>Reverse Comm Interf version</i> |
| /PPCG | <i>under construction</i> |
| LAXlib | <i>linear algebra library (int w ELPA)</i> |
| UtilXlib | <i>basic utilities (error,timinig,para)</i> |
| archive | <i>library archive (lapack source)</i> |
| clib | <i>c timing routine</i> |
| include | |
| install | <i>configure, makedeps</i> |
| Makefile | |
| configure | |



PPCG – Projected Preconditioned Conjugate Gradient

E. Vecharynski, C. Yang, J.E. Pask, *J. Comp.Phys.* **290**,73 (2015)

Algorithm 2: The projected preconditioned conjugate gradient (PPCG) algorithm.

Input: The matrix A , a preconditioner T , and a starting guess of the invariant subspace $X^{(0)} \in \mathbb{C}^{n \times k}$ associated with the k smallest eigenvalues of A ;

Output: An approximate invariant subspace $X \in \mathbb{C}^{n \times k}$ associated with the k smallest eigenvalues of A ;

1: $X \leftarrow \text{orth}(X^{(0)}); P \leftarrow []$;

2: **while** convergence not reached **do**

3: $W \leftarrow T(AX - X(X^*AX))$;

4: $W \leftarrow (I - XX^*)W$;

5: $P \leftarrow (I - XX^*)P$;

6: **for** $j = 1, \dots, k$ **do**

7: $S \leftarrow [x_j, w_j, p_j]$;

8: Find the smallest eigenpair $(\theta_{\min}, c_{\min})$ of $S^*ASc = \theta S^*Sc$, where $c^*S^*Sc = 1$;

9: $\alpha_j \leftarrow c_{\min}(1)$, $\beta_j \leftarrow c_{\min}(2)$; and $\gamma_j \leftarrow c_{\min}(3)$ ($\gamma_j = 0$ at the initial step);

10: $p_j \leftarrow \beta_j w_j + \gamma_j p_j$;

11: $x_j \leftarrow \alpha_j x_j + p_j$.

12: **end for**

13: $X \leftarrow \text{orth}(X)$;

14: If needed, perform the Rayleigh–Ritz procedure within $\text{span}(X)$;

15: **end while**

each band (or small group of bands) is updated by diagonalizing a small $3 \times \text{blksize} \times 3 \times \text{blksize}$ matrix built from the current X , the orthogonal residual and the orthogonal conjugate direction

- PPCG *work in progress*
- -memory friendly: bands are dealt with a small block at a time.
- -global calls to `h_psi` give opportunity for band parallelization (not working properly yet)
- -each block can be dealt with independently (parallelization)
- -most operations on arrays use efficient BLAS3 calls (DGEMM)

• routines

- `ppcg`, *real PPCG, cmplx version presently not available*
- `rotate_wfc_gamma`, *real initial diag (the same as CG)*
- `h_psi`, (`s_psi`) *generalized algorithm not available yet*
 - * `preconditioning`

Some recent work on an alternative iterative methods

A PARALLEL ORBITAL-UPDATING APPROACH FOR ELECTRONIC STRUCTURE CALCULATIONS *

XIAOYING DAI[†], XINGAO GONG[‡], AIHUI ZHOU[†] , AND JINWEI ZHU[†]

Abstract. In this paper, we propose an orbital iteration based parallel approach for electronic structure calculations. This approach is based on our understanding of the single-particle equations of independent particles that move in an effective potential. With this new approach, the solution of the single-particle equation is reduced to some solutions of independent linear algebraic systems and a small scale algebraic problem. It is demonstrated by our numerical experiments that this new approach is quite efficient for full-potential calculations for a class of molecular systems.

[arXiv:1405.0260v2 \[math.NA\] 20/11/2014](https://arxiv.org/abs/1405.0260v2)

A PARALLEL ORBITAL-UPDATING BASED OPTIMIZATION METHOD FOR ELECTRONIC STRUCTURE CALCULATIONS *

XIAOYING DAI[†], ZHUANG LIU[‡], XIN ZHANG[§], AND AIHUI ZHOU[¶]

Abstract. In this paper, we propose a parallel optimization method for electronic structure calculations based on a single orbital-updating approximation. It is shown by our numerical experiments that the method is efficient and reliable for atomic and molecular systems of large scale over supercomputers.

[arXiv:1510.07230v1 \[math.NA\] 25/10/2015](https://arxiv.org/abs/1510.07230v1)

ParO in a nutshell

ALGORITHM 1.1.

1. Given initial data $(\lambda_i^{(0)}, u_i^{(0)}) \in \mathbb{R} \times H_0^1(\Omega)$ with $(u_i^{(0)}, u_j^{(0)})_\Omega = \delta_{ij}$, ($i, j = 1, 2, \dots, N$), define \mathcal{T}_0 and V_0 , and let $n = 0$
2. Construct \mathcal{T}_{n+1} and V_{n+1} based on an adaptive procedure to $(\lambda_i^{(n)}, u_i^{(n)})$.
3. For $i = 1, 2, \dots, N$, find $u_i^{(n+1/2)} \in V_{n+1}$ satisfying

$$a(U^{(n)}; u_i^{(n+1/2)}, v) = \lambda_i^{(n)}(u_i^{(n)}, v) \quad \forall v \in V_{n+1}$$

in parallel.

4. Project to eigenspace: find $(\lambda^{(n+1)}, u^{(n+1)}) \in \mathbb{R} \times \tilde{V}_{n+1}$ satisfying $\|u^{(n+1)}\|_{0,\Omega} = 1$ and

$$a(U^{(n+1/2)}; u^{(n+1)}, v) = \lambda^{(n+1)}(u^{(n+1)}, v) \quad \forall v \in \tilde{V}_{n+1}$$

to obtain eigenpairs $(\lambda_i^{(n+1)}, u_i^{(n+1)})$ ($i = 1, 2, \dots, N$).

5. Let $n = n + 1$ and go to Step 2.

Here $\tilde{V}_{n+1} = \text{span} \{u_1^{(n+1/2)}, u_2^{(n+1/2)}, \dots, u_N^{(n+1/2)}\}$, $U^{(n)} = (u_1^{(n)}, u_2^{(n)}, \dots, u_N^{(n)})$, $U^{(n+1/2)} = (u_1^{(n+1/2)}, u_2^{(n+1/2)}, \dots, u_N^{(n+1/2)})$, and $a(\cdot; \cdot, \cdot)$ is the nonlinear variational form associated the Kohn-Sham equation defined in Section [2.2](#).

ParO as I understand it

- Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$
- Solve in parallel the *nbnd* linear systems

$$(H_{KS} + \lambda S)|\tilde{\phi}_i^{(n)}\rangle = (\varepsilon_i^{(n)} + \lambda)S|\phi_i^{(n)}\rangle$$

- Build the reduced Hamiltonian

$$\tilde{H}_{ij} = \langle \tilde{\phi}_i^{(n)} | H_{KS} | \tilde{\phi}_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \tilde{\phi}_i^{(n)} | S | \tilde{\phi}_j^{(n)} \rangle$$

- Diagonalize the small *nbnd* \times *nbnd* reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \quad \longrightarrow \quad \{|\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)}\}$$

- Repeat if needed in order to improve solution at fixed Hamiltonian

A variant of ParO method

• Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

• Solve in parallel the $nbnd$ linear systems

$$(H_{KS} + \lambda S)|\tilde{\phi}_i^{(n)}\rangle = (\varepsilon_i^{(n)} + \lambda)S|\phi_i^{(n)}\rangle$$

• Build the reduced Hamiltonian from both $|\tilde{\phi}_i^{(n)}\rangle$ & $|\phi_i^{(n)}\rangle$

$$\tilde{H}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | H_{KS} | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | S | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle$$

• Diagonalize the small $2nbnd \times 2nbnd$ reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \quad \longrightarrow \quad \{|\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)}\}$$

• Repeat if needed in order to improve solution at fixed Hamiltonian

A variant of ParO method (2)

• Given trial eigenpairs: $\{|\phi_i^{(n)}\rangle, \varepsilon_i^{(n)}\}$

• Solve in parallel the *nbnd* linear systems

$$\left(H_{KS} - \varepsilon_i^{(n)} S + \alpha S |\phi_i^{(n)}\rangle \langle \phi_i^{(n)}| S \right) |\tilde{\phi}_i^{(n)}\rangle = - (H_{KS} - \varepsilon_i^{(n)} S) |\phi_i^{(n)}\rangle$$

• Build the reduced Hamiltonian from both $|\tilde{\phi}_i^{(n)}\rangle$ & $|\phi_i^{(n)}\rangle$

$$\tilde{H}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | H_{KS} | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle, \quad \tilde{S}_{ij} = \langle \tilde{\phi}_i^{(n)} / \phi_i^{(n)} | S | \tilde{\phi}_j^{(n)} / \phi_j^{(n)} \rangle$$

• Diagonalize the small $2nbnd \times 2nbnd$ reduced Hamiltonian to get the new estimate for the eigenpairs

$$(\tilde{H} - \varepsilon \tilde{S})v = 0 \quad \longrightarrow \quad \{|\phi_i^{(n+1)}\rangle, \varepsilon_i^{(n+1)}\}$$

• Repeat if needed in order to improve solution at fixed Hamiltonian

- **Domain decomposition parallelization**
 - Basis set components are distributed
 - Memory is distributed
-
- **Band group parallelization**
 - Operations on dynamically defined band groups are distributed
 - Memory is NOT distributed
-
- **Parallel dense diagonalization**
 - A dedicated communicator is present (interface with ScaLapack and ELPA)